

Are jets rotating at the launching?

Noam Soker¹

¹ Department of Physics, Technion—Israel Institute of Technology, Haifa 32000 Israel;
 soker@physics.technion.ac.il.

Abstract. I argue that the Doppler shift asymmetries observed in some young stellar object (YSO) jets result from the interaction of the jets with the circumstellar gas, rather than from jets' rotation. The jets do rotate, but at a velocity much below claimed values. During the meeting I carefully examined new claims, and found problems with the claimed jets' rotation. I will challenge any future observation that will claim to detect jet rotation in YSOs that requires the jets (and not a wind) to be launched from radii much larger than the accreting stellar radius. I conclude that the most likely jets' launching mechanism involves a very efficient dynamo in the inner part of the accretion disk, with a jets launching mechanism that is similar to solar flares (coronal mass ejection).

1. Introduction

In previous papers (Soker 2005, 2007a) I proposed that the interaction of the jets with a twisted-tilted (wrapped) accretion disk can form the asymmetry in the jets' line of sight velocity profiles as observed in some YSOs (e.g. Bacciotti et al. 2002). The claim that the observations of asymmetric Doppler shifts do not support jet rotation in YSOs was strengthened by the numerical simulations of Cerqueira et al. (2006). They assumed a precessing jet whose ejection velocity changes periodically with a period equals to the precession period. Practically, the dependance of the jets expansion velocity on direction around the symmetry axis leads to the same effect as the model of Soker (2005). Whereas in Soker (2005) the physical process behind this velocity profile is an interaction with the material in the jets surroundings, Cerqueira et al. (2006) give no justification for the periodic variation of the jets ejection speed. As far as comparison with observation is considered, it is hard to distinguish between the model of jet interaction with its surrounding (Soker 2005), and the periodic jets speed of Cerqueira et al. (2006).

2. Problems with claimed jets' rotation

To demonstrate the problems with the argued jet rotation, I will examine two new claims.

After the publication of my earlier papers Zapata et al. (2009) argued for a rotating molecular jet in Ori-S6. I find four problems with this case. More detail can be found in my presentation at the meeting:

<http://iaujd-outflows.blogspot.com/2008/10/scientific-program.html>

(1) In some regions the red and blue shifted components overlap. This is against expectation if the red-blue shifted components are due to jets' rotation. (2) In some regions the blue and red shifted components are disconnected. As each jet is one entity, this is against expectation if the Doppler shifts are due to jets' rotation. (3) Using the rotation interpretation at the edge of $30''$ across the disk, gives a jets foot-point of 300 AU. This is larger than the size of the accretion disk given by the same authors for this object.

(4) The ring that supposedly feeds the accretion disk and the jets, rotates in opposite sense to that of the claimed jets' rotation. As Zapata et al. write: "The sense of rotation of the circumbinary ring is nearly opposite to that of jet and outflow, and the jet leaves the system under an angle of 45° with the ring plane." I note that a tilted jet can lead to the asymmetric red-blue shift, as in the model I proposed in 2005.

During the meeting, I was challenged to account for a very recent claim of a possible jets' rotation in HH 211 (Lee et al. 2009, 2007). I find two problems with the tentative claimed jets' rotation (I elaborate on these points in the appendix in the astro-ph version of this paper). (1) The blue and red components exchange sides. Namely, the velocity plots do not give a clear sense of asymmetry, and hence no unique sense of rotation. The same effect is seen in the velocity maps of HH 212 (Lee et al. 2008). (2) The accretion disk cannot supply the required energy and angular momentum if the rotation is real.

My conclusion is that these types of observations give peaks in emission that show different Doppler shifts. By pure fluctuations, these might mimic rotation in some places. In some cases the sense of the fluctuations will give rotation in the same sense as that of the accretion disk. In other cases the sense will be in an opposite direction to that of the disk, and in some cases just zero rotation will be deduced. The inferred rotation is due to fluctuations that by chance can mimic rotation.

3. The launching mechanism

The talks and discussions during the meeting strengthened my view that the launching mechanism involves reconnection of magnetic field lines. Reconnection can occur between the stellar and the disk magnetic fields (e.g., de Gouveia dal Pino & Lazarian 2005; de Gouveia dal Pino et al. 2009), or reconnection of the disk magnetic field (Soker 2007b). Laor & Behar (2008) show that the ratio of radio luminosity to X-ray luminosity has similar values in magnetically active stars and in many accreting objects, up to radio quiet quasars. Based on this correlation I prefer the following conclusion (Soker 2007b; Soker & Vrtillek 2009): There is a very efficient dynamo in the inner part of the accretion disk, with a jets launching mechanism that is similar to solar flares (coronal mass ejection).

References

- Bacciotti, F., Ray, T. P., Mundt, R., Eisloffel, J., & Solf, Jo. *ApJ*, 576, 222 (B2002)
- Cerqueira, A. H., Velazquez, P. F., Raga, A. C., Vasconcelos, M. J., & de Colle, F. 2006, *A&A*, 448, 231
- de Gouveia dal Pino, E. M., & Lazarian, A. 2005, *A&A*, 441, 845
- de Gouveia Dal Pino, E. M., Piovezan, P., Kadowaki, L., Kowal, G., & Lazarian, A., this proceedings (IUA JD 7 at the XXVIIth IAU General Assembly)
- Laor, A., & Behar, E. 2008, *MNRAS*, 390, 847
- Lee, C.-F., et al. 2007, *ApJ*, 670, 1188
- Lee, C.-F., et al. 2008, *ApJ*, 685, 1026
- Lee, C.-F., et al. 2009. *ApJ*, 699, 1584
- Soker, N. 2005, *A&A*, 435, 125
- Soker, N. 2007a, astro-ph/0703474
- Soker, N. 2007b, IAUS, 243, 195 (arXiv:0706.4241).
- Soker, N., & Vrtillek, S. D. 2009, arXiv:0904.0681
- Zapata, L. A., Schmid-Burgk, J., Muders, D., Schilke, P., Menten, K., & Guesten, R. 2009, *A&A* in press

4. Appendix (astro-ph version only): Do the jets in HH211 rotate?

I took the challenge raised during my talk, and looked at the papers by Lee et al. on HH211 (Lee et al. 2007, 2009). There are two problematic points.

4.1. Velocity Maps

In figure 11 of Lee et al. (2009) they show a 1 km/sec velocity map. They claim the red component peaks on one side, while the blue on the other. However, it seems that the blue and red components exchange sides, as marked on the figure below. From what I can tell from other figures in their second paper, the velocity plots do not give a clear sense of asymmetry. The same feature appears in the velocity maps of HH 212 presented by Lee et al. (2008). In their figs. 4 and 6 the sign of the velocity gradient changes between different locations. In their fig. 1b the red and blue components exchange sides. Here as well, fluctuations seems to cause these variations.

4.2. The foot-point

In the first paper they derive the two-sided mass-loss rate of the jets

$$\dot{M}_{2j} = (0.7 - 2.8) \times 10^{-6} M_{\odot} \text{ yr}^{-1}. \quad (4.1)$$

For the accretion rate they give the following values for the first and second paper, respectively

$$\dot{M}_{\text{acc-f}} = 8 \times 10^{-6} M_{\odot} \text{ yr}^{-1}; \quad \dot{M}_{\text{acc-s}} = (2.5 - 5) \times 10^{-6} M_{\odot} \text{ yr}^{-1}. \quad (4.2)$$

Taken all these values, the ejection in the jets to accretion ratio is

$$\eta_f \equiv \frac{\dot{M}_{\text{acc-f}}}{\dot{M}_{2j}} \simeq 0.09 - 0.36; \quad \eta_s \equiv \frac{\dot{M}_{\text{acc-s}}}{\dot{M}_{2j}} \simeq 0.2 - 0.6, \quad (4.3)$$

by the first and second paper, respectively.

A jet velocity of $v_j = 170 \text{ km s}^{-1}$ from their star of mass $M = 0.05 M_{\odot}$, and the above ratio η , require the foot point to be at radius r_0 determined from energy conservation

$$\frac{1}{2} \frac{GM}{r_0} \dot{M}_{\text{acc}} = \frac{1}{2} \dot{M}_{2j} v_j^2. \quad (4.4)$$

This gives

$$r_0 = \frac{GM}{v_j^2} \eta^{-1} = 3.3 \left(\frac{\eta}{0.1} \right)^{-1} R_{\odot}. \quad (4.5)$$

The specific angular momentum that the accreted mass can supply to the jet is

$$j_m = \frac{\sqrt{GM r_0}}{\eta} = 8.3 \left(\frac{\eta}{0.1} \right)^{-1} \left(\frac{r_0}{3.3 \text{ AU}} \right)^{-1} = 8.3 \left(\frac{\eta}{0.1} \right)^{-2} \text{ km s}^{-1} \text{ AU}. \quad (4.6)$$

The claimed observed value is $j_m = 5 \text{ km s}^{-1} \text{ AU}$. Namely, the values of η is constraint to be $\eta < 0.15$. It cannot be as large as 0.3 as it is according to Lee et al. (2009). As I do not expect that the accreted mass will transfer angular momentum and energy at 100% efficiency to the ejected gas (namely, its rotation speed will not drop to zero), the constraint on η is stronger even, $\eta < 0.1$.

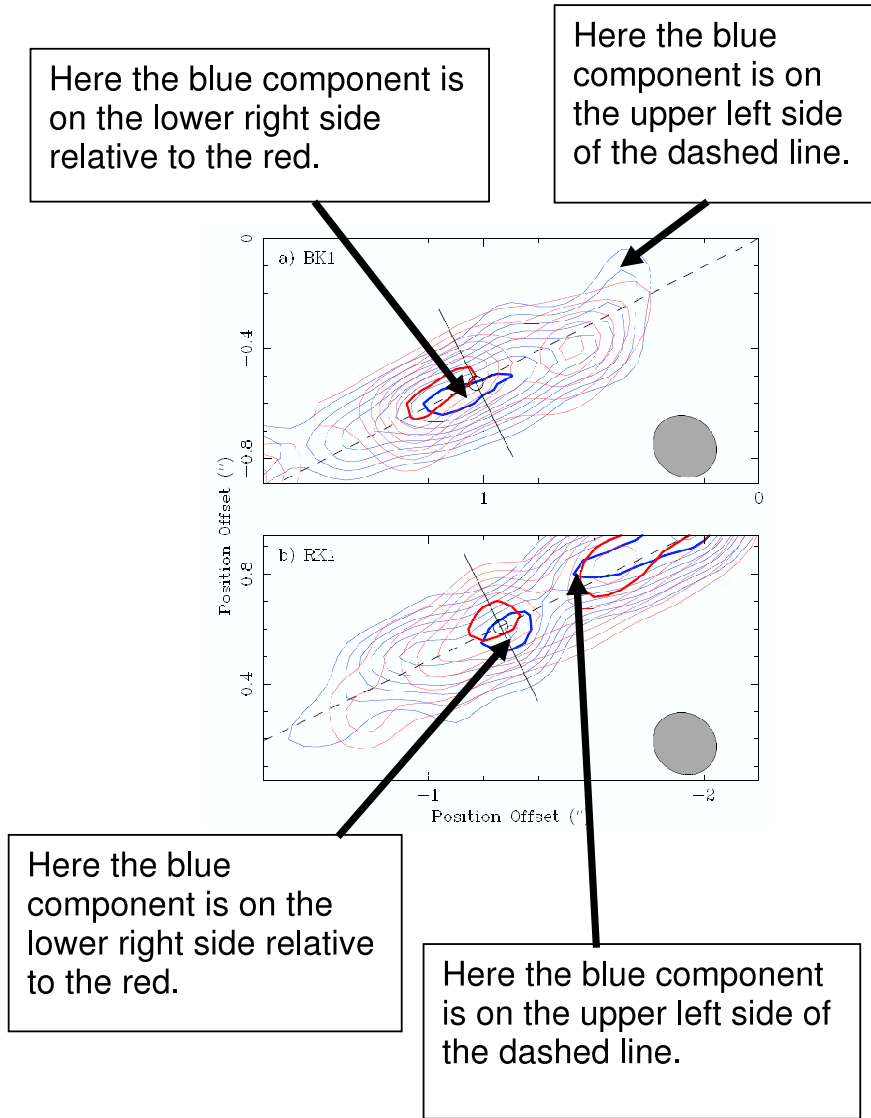


Figure 1. The velocity contour map from figure 11 of Lee et al. (2009). Marked are the changes in the sense of rotation.